

Transversity: theory vs experiments

- Why transversity? [h_1 , $\Delta_T q$, δq]
Experimental access to transversity
- h_1 in $p^\dagger p^\dagger$ interactions [$h_1 \otimes h_1$]
- New chiral-odd fragmentation functions
- h_1 in $pp^\dagger \rightarrow \pi X$ or $pp^\dagger \rightarrow \Lambda^\dagger X$
- h_1 in SIDIS processes:
 - $l p^\dagger \rightarrow l \pi X$
 - $l p^\dagger \rightarrow l \pi \pi X$
 - $l p^\dagger \rightarrow l V^\dagger X$
 - $l p^\dagger \rightarrow l \Lambda^\dagger X$
- Access to new frag. functs.
 - $e^+e^- \rightarrow \pi^+ \pi^- X$
 - $e^+e^- \rightarrow \pi \pi \pi \pi X$
- Conclusions

Trieste, 19/2/2002
M. ANSELMINO

Transversity, h_1 or $\Delta_T q$ or δq

q , Δq and h_1 (or $\Delta_T q$, δq):

fundamental leading-twist quark distributions [equally important]

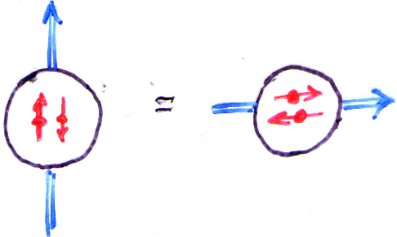
q = quark distribution ($= \int q/p$)

Δq = helicity distribution

$$= q_+^+ - q_-^+ \quad (= \int_{q_+/P_+} - \int_{q_-/P_+})$$

$\Delta_T q$ = transverse polarization distribution

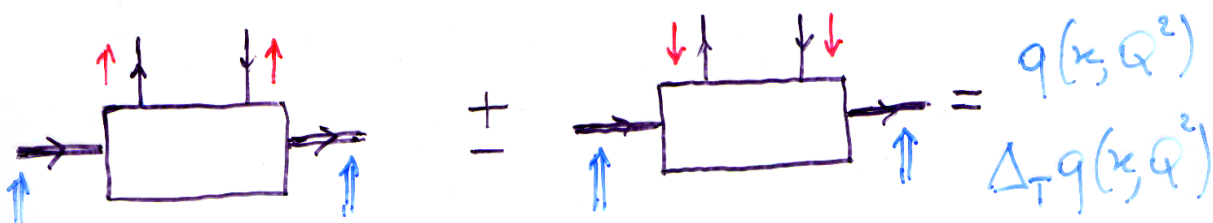
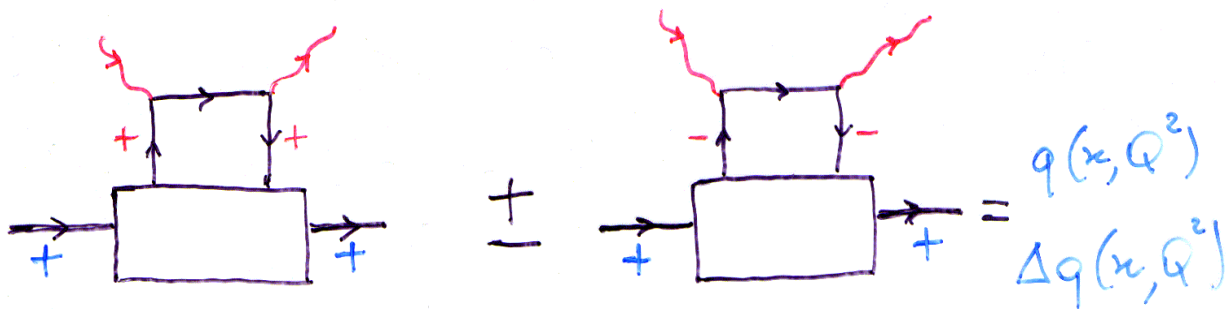
$$= q_{\uparrow}^{\uparrow} - q_{\downarrow}^{\uparrow}$$

$\Delta_T q = \Delta q$ only at rest 

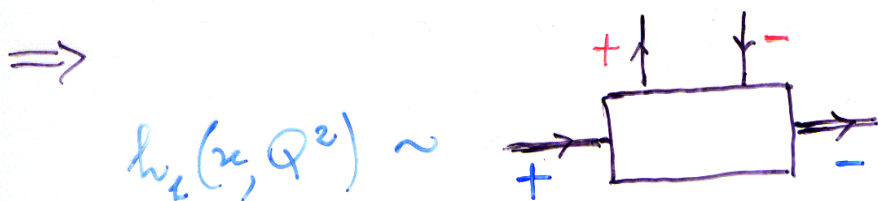
$$\Delta q \sim \bar{q} \gamma^M \gamma_5 q \quad (\text{chiral-even})$$

$$\Delta_T q \sim \bar{q} \sigma^{\mu\nu} \gamma_5 q \quad (\text{chiral-odd})$$

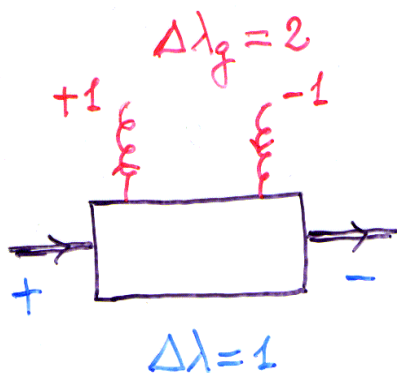
q and Δq can be measured in DIS



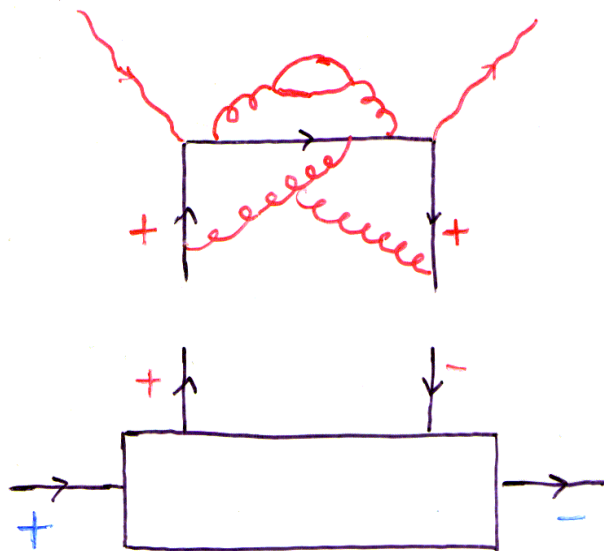
In helicity basis, $\uparrow, \downarrow = \frac{1}{\sqrt{2}} (|+\rangle \pm i|-\rangle)$



\Rightarrow no gluon contribution to h_2



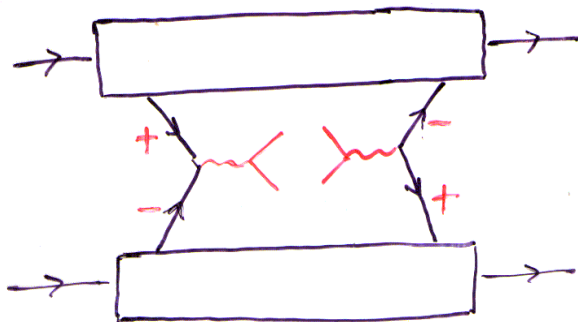
(simple Q^2 evolution)



h_1 decouples from QCD, SM dynamics

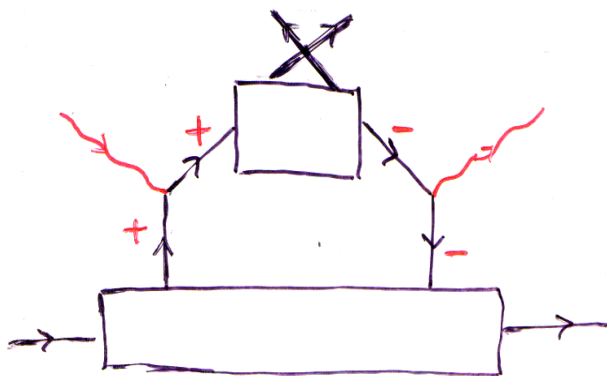
$\Rightarrow h_1$ must couple to another chiral-odd function. For example:

Drell-Yan processes ($pp \rightarrow \mu^+ \mu^- X$)



$\sim h_1 \otimes h_1$

SIDIS ($lp \rightarrow lh X$)



$\sim h_1 \otimes \tilde{D}_{G-O}$

h_L from $p^\uparrow p^\uparrow$ processes

$$\text{Measure } A_{TT} = \frac{d\sigma^{\uparrow\uparrow} - d\sigma^{\uparrow\downarrow}}{d\sigma^{\uparrow\uparrow} + d\sigma^{\uparrow\downarrow}} \equiv \frac{\Delta_T \sigma}{\sigma}$$

- $p^\uparrow p^\uparrow \rightarrow \mu^+ \mu^- X$

$$\Delta_T \sigma^{\mu\mu} \sim \sum_q h_{Lq} \otimes h_{L\bar{q}} \otimes \Delta_T \hat{\sigma}^{\uparrow}$$

↙ very small

- $p^\uparrow p^\uparrow \rightarrow \pi^+ \pi^- X$

$$\Delta_T \sigma^{\pi\pi} \sim \sum_{abcd} h_{La} \otimes h_{Lb} \otimes \Delta_T \hat{\sigma} \otimes D_{\pi/c} \otimes D_{\pi/d}$$

no gluon contribution $\Rightarrow A_{TT}$ very small

$$A_{TT} = \frac{\Delta_T \sigma^{\pi^+ \pi^+} + \Delta_T \sigma^{\pi^- \pi^-} - \Delta_T \sigma^{\pi^+ \pi^-} - \Delta_T \sigma^{\pi^- \pi^+}}{\sigma^{\pi^+ \pi^+} + \sigma^{\pi^- \pi^-} - \sigma^{\pi^+ \pi^-} - \sigma^{\pi^- \pi^+}}$$

large, but huge statistical errors
(de Florian, Stratmann, Vogelsang)

- $p^\uparrow p^\uparrow \rightarrow 2 \text{ jets} + X$

A_{TT} small ($\sim 1\%$) but very good statistic at RHIC

New chiral-odd fragmentation facts.

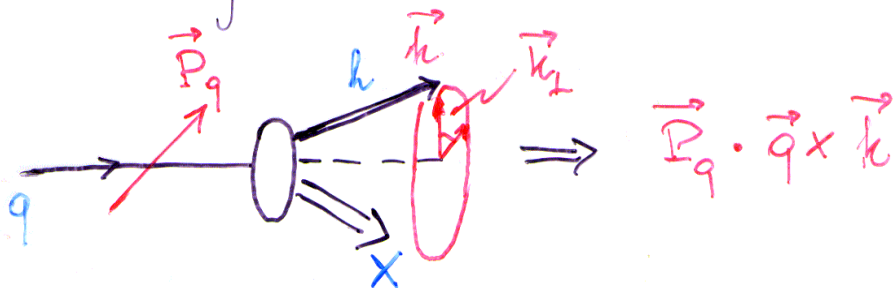
Access h_L via chiral-odd f.f.

Which ones?

- Analogue of h_L for fragmentation

$$\Rightarrow \Delta_T D = D_{q\uparrow}^{h\uparrow} - D_{q\downarrow}^{h\uparrow}$$

- Collins function



$$\Delta^N D_{h/q\uparrow} = D_{h/q\uparrow}(z, \vec{k}_+) - D_{h/q\downarrow}(z, \vec{k}_-) \quad (\sim H_L^+)$$

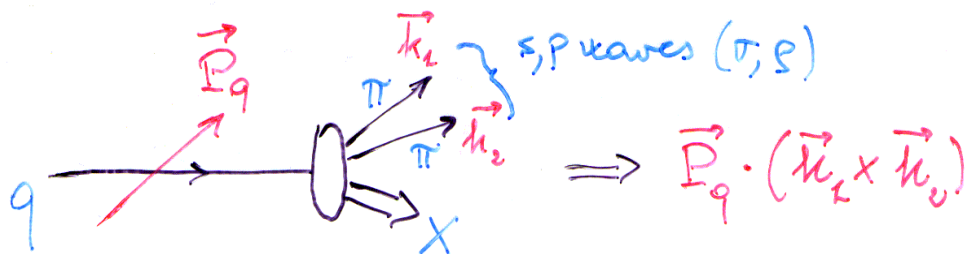
$$\sim \text{Im} \sum_{\lambda_h} \sum_{\sigma_x, \lambda_x} D_{\lambda_h \lambda_x; +} D_{\lambda_h \lambda_x; -}^*$$

(needs interferences between frag. amplitudes)

$$D_{h/q}(z, \vec{k}_\pm) = \frac{1}{2} \sum_{\lambda_q \lambda_h} \sum_{\sigma_x, \lambda_x} |D_{\lambda_h \lambda_x; \lambda_q}|^2$$

[H_L^+ , Bacchetta, Kumer, Metz, Mulders]

- Interference frag. functions
(Jaffe, Ji, Tang; Artru, Collins)



Int. between s,p $\pi\pi$ states \Rightarrow new f.f.

$\delta\hat{q}_I$ [for a model see Radici, Jakob, Bianconi]

- For production of spin \pm (V) particles
[M.A., Boglione, Haiman, Murgia]

$$D_{4,0}^{+,-} \equiv \sum_{\lambda, \lambda'} D_{\lambda \times \lambda, +} D_{\lambda \times 0, -}^*$$

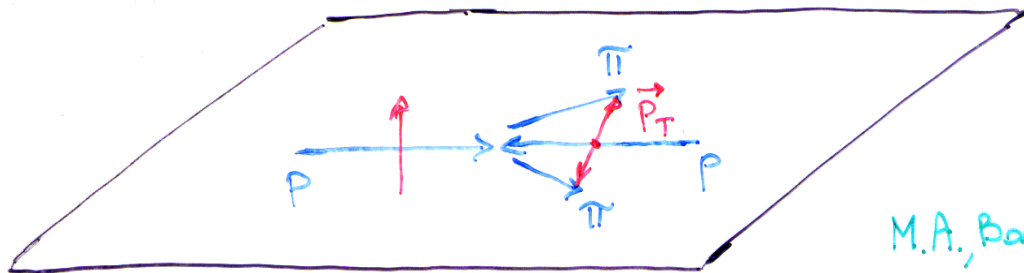
contributes to $f_{10}(V)$ of vector particles
produced in SIDIS

[See also Bacchetta, Mulders]

- Transversal handedness
[Efremov et al.]

Which processes?

$$- A_N \text{ in } p^{\uparrow} p \rightarrow \pi^{\uparrow} X = \frac{d\sigma^{\uparrow} - d\sigma^{\downarrow}}{d\sigma^{\uparrow} + d\sigma^{\downarrow}}$$



M.A. Baglione
Leader, Murgia

left-right asymmetry

$$d\sigma^{\uparrow} - d\sigma^{\downarrow} \sim \sum_{abc} h_{1a} \otimes f_{b/p} \otimes \Delta\sigma \otimes \Delta^N D_{\pi/c}^{\uparrow}$$

\downarrow trans. \downarrow Collins f.

$$- P_T^{\Lambda} \text{ in } p p^{\uparrow} \rightarrow \Lambda^{\uparrow} X = \frac{d\sigma^{\Lambda^{\uparrow}} - d\sigma^{\Lambda^{\downarrow}}}{d\sigma^{\Lambda^{\uparrow}} + d\sigma^{\Lambda^{\downarrow}}}$$

[Λ in current fragm. region]

$$P_T^{\Lambda} \sim \sum_{abc} f_{a/p} \otimes h_{1b} \otimes \Delta\sigma \otimes \Delta_T D_{\Lambda/c}$$

\downarrow trans. \downarrow transverse fragmentation

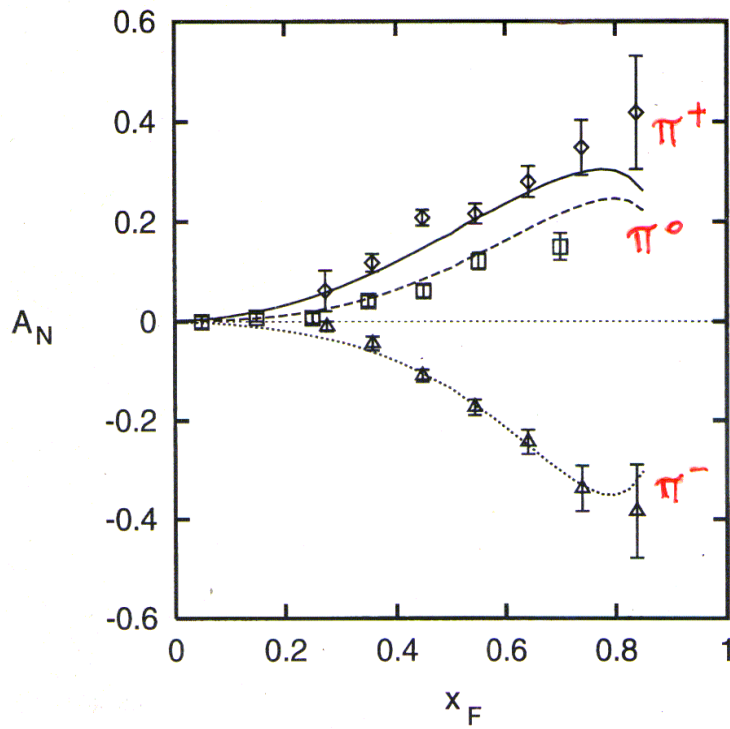


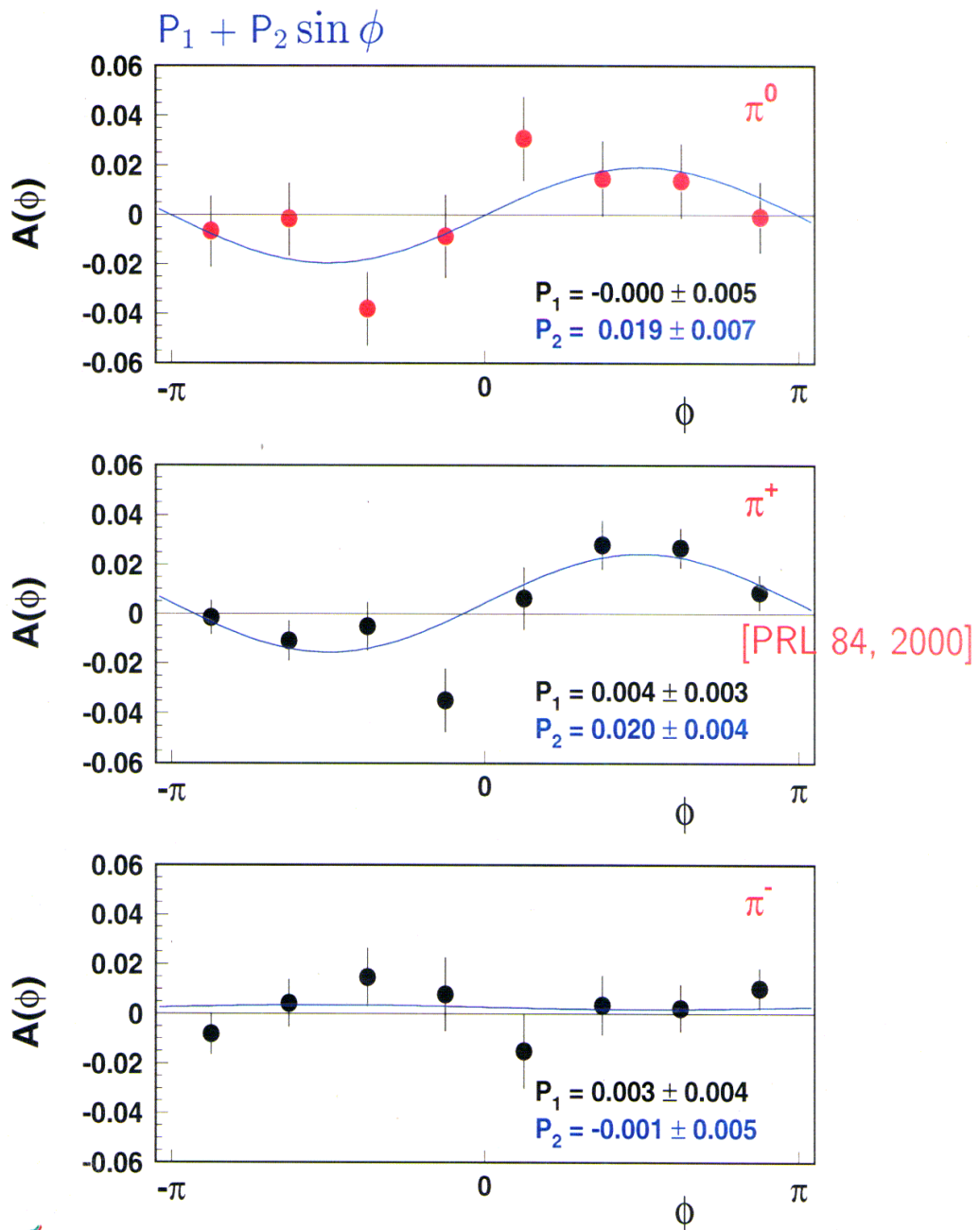
Figure 2: Single spin asymmetry for pion production, $p^\uparrow p \rightarrow \pi X$. The data points are the E704 experimental single spin asymmetry for π^+ (diamonds), π^0 (squares) and π^- (triangles) [3]. The solid line is our best fit for π^+ , the dashed line for π^0 and the dotted line for π^- , obtained under the assumption of Collins effect only.

$$\Delta^N D_{\pi/q^+} \neq 0$$

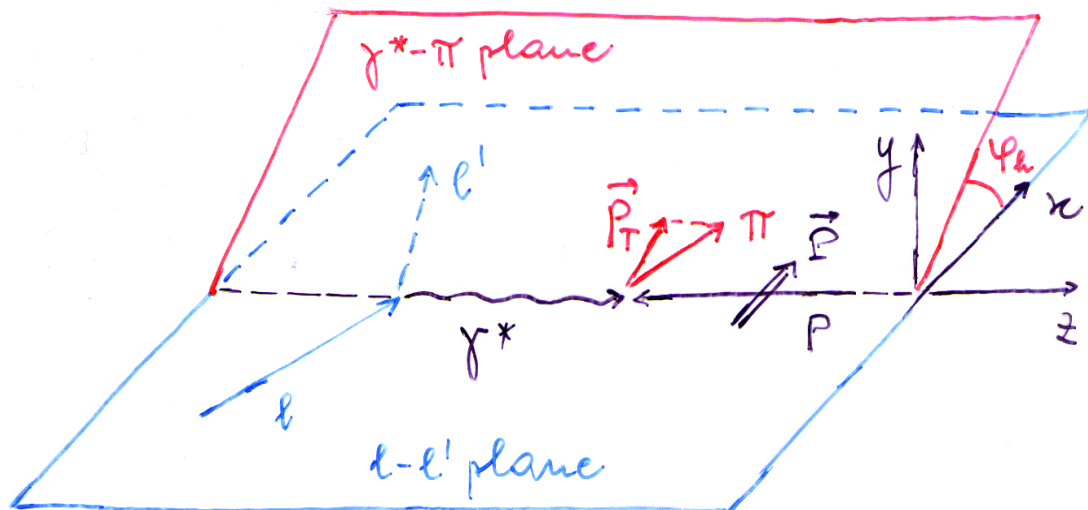
$lp^{\uparrow} \rightarrow l\pi X$

single target-spin azimuthal asymmetry

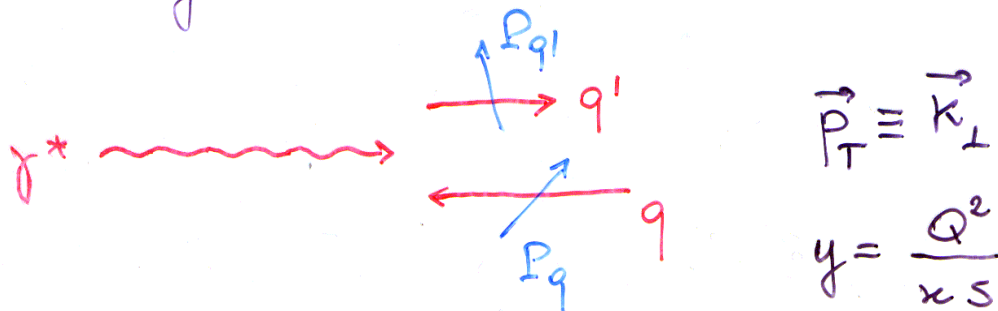
$$A(\phi) = \frac{1}{\langle P_t \rangle} \cdot \frac{N^+(\phi) - N^-(\phi)}{N^+(\phi) + N^-(\phi)}$$



$l p^\uparrow \rightarrow l \pi X$ cross-section in $\gamma^* - p$ c.m.



Elementary interaction:

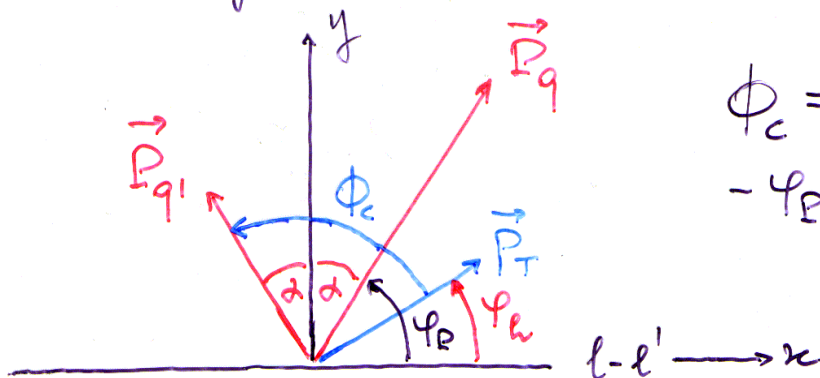


$$\vec{P}_T \equiv \vec{k}_T$$

$$y = \frac{Q^2}{x s}$$

$$\begin{cases} (P_{q'})_x = -D_{NN} (P_q)_x \\ (P_{q'})_y = D_{NN} (P_q)_y \end{cases}$$

$$D_{NN} = \frac{2(1-y)}{1+(1-y)^2}$$



$$\phi_c = \pi + \theta - \phi_p - \phi_h$$

$$b p^\uparrow \rightarrow b \pi X \quad A_N = \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow}$$

At leading twist:

$$A_N^b = \frac{\sum_q e_q^2 h_{2q} \Delta^N D_{h/q}^\uparrow}{\sum_q e_q^2 q \cdot 2 D_{h/q}} \frac{2(1-y)}{1+(1-y)^2} P \sin \phi_c$$

Using isospin and charge conjugation invariance ($D_{\pi^+/u} = D_{\pi^-/d}$ etc.) and neglecting non valence contributions ($D_{\pi^+/d} \approx 0 \dots$)

$$\Rightarrow A_N^{\pi^i} = \frac{h_i(x)}{f_i(x)} \frac{\Delta^N D_{\pi/q}^\uparrow(z, P_T)}{2 D_{\pi/q}(z, P_T)} \frac{2(1-y)}{1+(1-y)^2} P \sin \phi_c$$

$$i = + \quad h_+ = 4h_{cu} \quad f_+ = 4u + \bar{d}$$

$$i = - \quad h_- = h_{cd} \quad f_- = d + 4\bar{u}$$

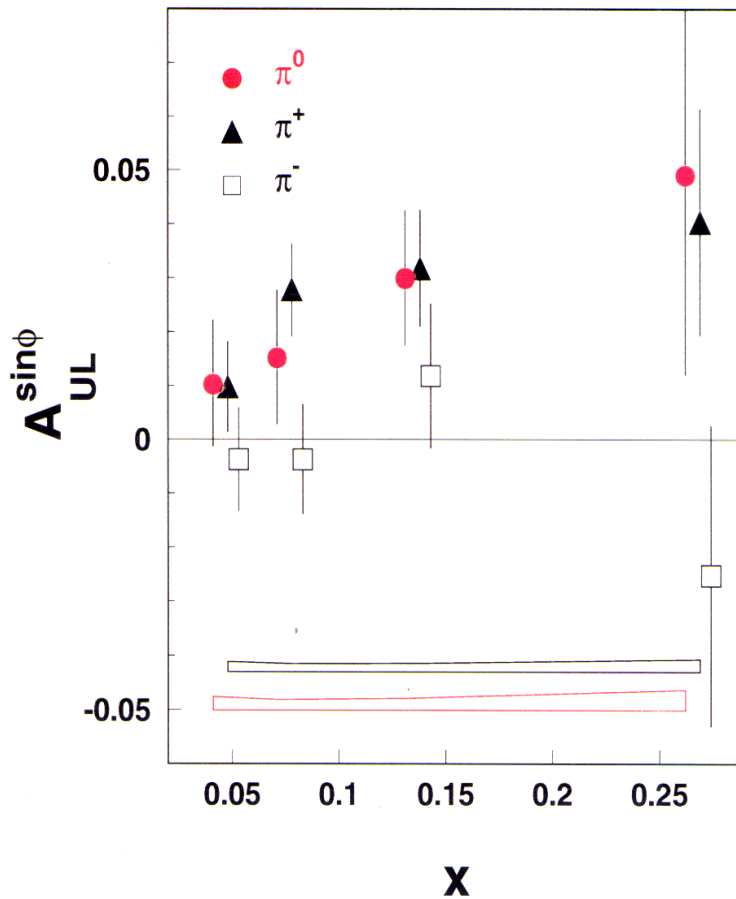
$$i = 0 \quad h_0 = 4h_{cu} + h_{cd} \quad f_0 = 4u + d + 4\bar{u} + \bar{d}$$

One expects:

$$A_N^{\pi^+} \approx A_N^{\pi^0} \quad \text{at large } x$$

$A_N^{\pi^i}$ increases with x and P_T

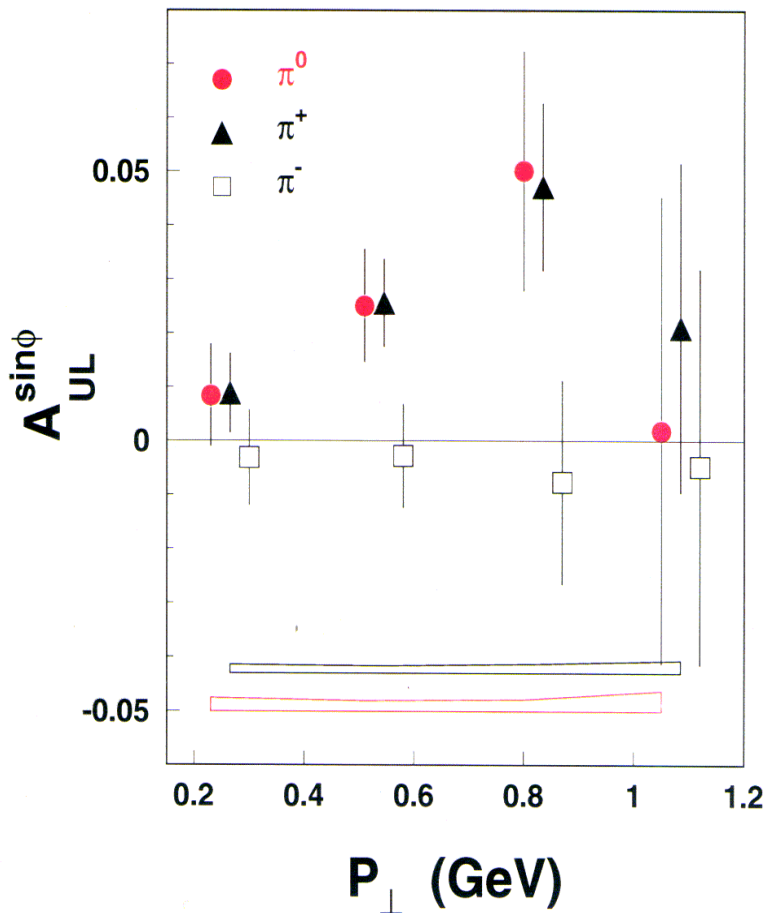
SSA $A_{UL}^{\sin\phi}$: x dependence



similar behaviour of π^0 and π^+

behaviour consistent with increase suggests that single-spin asymmetries are associated with valence quark contributions

SSA $A_{UL}^{\sin\phi}$: P_{\perp} dependence



related to the dominant kinematic role of the intrinsic transverse momentum of the quark, if P_{\perp} remains below the typical hadronic mass of ~ 1 GeV

Compare

$$A_N^{\pi^+} = \frac{4 l_{2u}}{4u + \bar{d}} \frac{\Delta^N D_{\pi^+/q}}{2 D_{\pi^+/q}} \frac{2(1-y)}{1+(1-y)^2} \sin \phi_c$$

with SMC data (not suppressed by $\frac{1}{Q}$ factors)

$$A_N^{\pi^+} \approx -(0.10 \pm 0.06) \sin \phi_c$$

Use Soffer's bound on l_{2u}

$$|l_{2q}| \leq \frac{1}{2} (q + \Delta q)$$

q and Δq are known

\Rightarrow obtain lower bound on $\frac{|\Delta^N D|}{D}$:

$$\frac{|\Delta^N D_{\pi^+/q}|}{D_{\pi^+/q}} \geq 0.24 \pm 0.15 \quad \text{large!}$$

Similar results from HERMES data
(but higher twist might be important)

[M.A., F. Murgia]

- A_N in $l p^\uparrow \rightarrow l \pi \pi X$ (or $p p^\uparrow \rightarrow \pi \pi X$)

$$d\sigma^\uparrow - d\sigma^\downarrow \sim \sin(\delta_0 - \delta_2) \sum_q \frac{e_q^2}{q} h_{1q}(x) \delta \hat{q}_I(z)$$

- $P_{10}(V)$ in $l p^\uparrow \rightarrow l V X$ ~ elem. int.

$$\text{Re } P_{10}(V) d\sigma \sim \sum_q h_{1q}(x) \hat{M} \hat{M}^*(y) \text{Im } D_{1,0}^{+,-}$$

simple model:

$$\begin{aligned} \text{Im } D_{1,0}^{+,-} &\approx [D_{V_2/q_+} D_{V_0/q_-}]^{\frac{1}{2}} \\ &\approx \frac{\sqrt{2}}{3} D_{V/q} \end{aligned}$$

- P_T^\uparrow in $l p^\uparrow \rightarrow l \Lambda^\uparrow X$

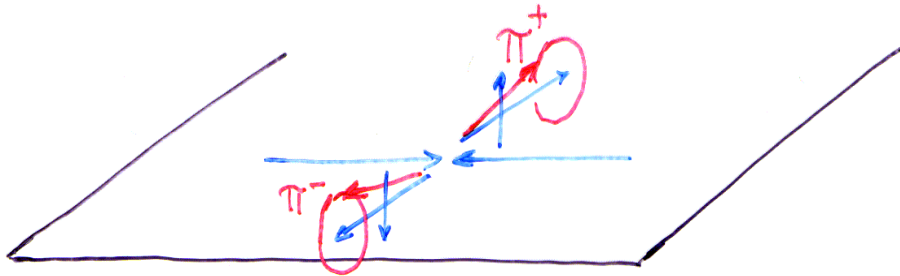
$$P_T^\uparrow = \frac{\sum_q \frac{e_q^2}{q} h_{1q}(x) \Delta_T D_{\Lambda/q}(z)}{\sum_q \frac{e_q^2}{q} q(x) D_{\Lambda/q}(z)} \frac{2(1-y)}{1+(1-y)^2}$$

~ separate x, y, z dependences

Λ 's in current frag. region

Direct access to new f.f.?

- $e^+e^- \rightarrow \pi^+\pi^- X$



\Rightarrow Dependence on angle ϕ between the two planes ($e^+e^-\pi^+$) and ($e^+e^-\pi^-$)

$$d\sigma^{\pi^+\pi^-} \sim (1 + \cos 2\phi) \Delta^N D_{\pi^+/q^+} \Delta^N D_{\pi^-/\bar{q}^-}$$

Boer, Jakob, Mulders (BELLE-KEK)

- $e^+e^- \rightarrow (\pi\pi)(\pi\pi) X$

\Rightarrow ang. correlation $\sim \delta \hat{q}_I \delta \hat{q}_I$

- $e^+e^- \rightarrow \Lambda^+ \bar{\Lambda}^+ X$

\Rightarrow information on $\Delta_T D_{\Lambda/q} \Delta_T D_{\bar{\Lambda}/\bar{q}}$

Conclusions

- $\Delta_T q$ has same status and rights as q and Δq
- Much harder to measure, as it is chiral-odd
- Couple $\Delta_T q$ with another chiral-odd function \Rightarrow
 $\Delta_T q \otimes \Delta_T q$, $\Delta_T q \otimes \Delta^N D$,
 $\Delta_T q \otimes \Delta_T D$, $\Delta_T q \otimes \hat{\sigma} q_I$, ...
- Look for separate variable dependences, separate information on new functions, combined experimental information, models for new functions, ...

[Review paper by Barone, Drago, Ratcliffe]